

Large time behavior for the heat equation on Carnot groups

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Abstract. We first generalize a decomposition of functions on Carnot groups as linear combinations of the Dirac delta and some of its derivatives, where the weights are the moments of the function. We then use the decomposition to describe the large time behavior of solutions of the hypoelliptic heat equation on Carnot groups. The solution is decomposed as a weighted sum of the hypoelliptic fundamental kernel and its derivatives, the coefficients being the moments of the initial datum.

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1. Introduction

The study of Carnot groups and of PDEs on these spaces have drawn an increasing attention for several reasons. First, Carnot groups are topologically extremely simple, since they are isomorphic to \mathbb{R}^N , but in the same time their metric structure is different, since they are naturally endowed with Carnot–Caratheodory metrics. Second, several tangent spaces to Carnot–Caratheodory metric spaces are diffeomorphic to Carnot groups. Third, and more connected to PDEs, Carnot groups are the easiest examples of spaces where hypoelliptic equations are naturally defined (see [4] and references therein). Thus, they are the most natural examples in which one can study relations between the solutions of the hypoelliptic PDEs and the Carnot–Caratheodory metrics.

With this goal, we first generalize a decomposition of functions first stated in [7] for functions on \mathbb{R}^N . In that article, the authors find conditions to write functions as

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$$f = a_0\delta_0 + \sum_{i=1}^N a_i\partial_i\delta_0 + \sum_{i,j=1}^N a_{ij}\partial_{ij}\delta_0 + \dots, \quad (1)$$

where δ_0 is the Dirac delta at 0 and $\delta_i, \delta_{ij}, \dots$ are standard derivatives in \mathbb{R}^N . The coefficients a_0, a_i, a_{ij}, \dots are called the moments. Our goal is to extend such kind of formula to Carnot groups. In this case, we have to replace \mathbb{R}^N derivatives with derivatives that are “adapted” to the Carnot group. The natural choice is to use derivative operators that are invariant with respect to the operation of the Carnot group. We will recall all the necessary definitions and result for Carnot groups in Sect. 2.

One of the interest of (1) is that it is useful to study the heat equation on Carnot groups. In this case, the standard Laplacian $\Delta = \sum_{i=1}^N \delta_i^2$ is replaced by the sub-Riemannian Laplacian Δ_{sr} , that is an hypoelliptic operator. Also in this case, all necessary definitions and result are recalled in Sect. 2. In particular, we study the initial value problem

$$\begin{cases} \partial_t f = \Delta_{sr} f \\ f(t=0) = f_0, \end{cases} \quad (2)$$

whose solution is given by $f(t) = f_0 * P_t$, where the convolution $*$ is intrinsically defined on the Carnot group G by its Lie group structure (see Definition 2.10) and P_t is the fundamental solution of the hypoelliptic heat equation with initial datum δ_0 . More precisely, we study the large time behavior of solutions of (2) and present a complete asymptotic expansion using the fundamental solution and its derivatives as basis functions and the moments of the initial datum as coefficients. We show that, due to anisotropy, only a subset of the moments needs to be used.

The structure of the paper is the following. In Sect. 2, we recall the definition of Carnot groups, of the metrics defined on them, the convolution and some useful properties of the fundamental solutions for the heat equation. Section 3 is devoted to the main result of this article, the decomposition of a function with respect to its moments. Section 4 is devoted to the study of the initial value problem (2), with a description of the large time behaviour in terms of the moments of the initial datum. We finally apply these results to the case of the Heisenberg group.

2. Carnot groups and hypoelliptic heat equations

We define here Carnot groups and recall some of their fundamental properties. For more details, see [4, 9].

Let \mathfrak{L} be a nilpotent Lie algebra that, by definition, is a vector space, thus isomorphic to \mathbb{R}^N for some $N \geq 1$. A Lie algebra is endowed with a Lie bracket product $[\cdot, \cdot]$. One can always assume that the Lie algebra is an algebra of matrices (Ado theorem, see e.g. [3]), and that the Lie bracket is $[A, B] = AB - BA$. We recall that the Lie algebra is nilpotent where there exists a $n > 0$ such that $[A_1, [A_2, [A_3, \dots [A_{n-1}, A_n] \dots]] = 0$ for any choice of $A_1, \dots, A_n \in \mathfrak{L}$. In terms of matrices, a nilpotent algebra can be always

identified with a subalgebra of the algebra of upper triangular matrices of a given dimension with zero diagonal.

Let G be the exponential of \mathfrak{L} . Again, in terms of matrices, one can consider the exponential map $\exp : \mathfrak{L} \rightarrow G$ to be the standard matrix exponential. G is a group with respect to the standard matrix product, and it is indeed a Lie group. For more details, see e.g. [3]. Since \mathfrak{L} is nilpotent, the exponential map $\exp : \mathfrak{L} \rightarrow G$ is a diffeomorphism. We fix a basis¹ $\{X^1, \dots, X^N\}$ on \mathfrak{L} . All along the paper, the notation $x = (x_1, \dots, x_N)$ can denote both $x = \sum x_i X^i \in \mathfrak{L}$ or $x = \exp(\sum x_i X^i) \in G$, depending on the context. Also recall that the Lie algebra of a Lie group is identified with the space of invariant differential operators on the Lie group, i.e. of vector fields X such that $X(g) = (L_g)_* X(e)$, where L_g is the left translation $h \rightarrow gh$ and $(L_g)_*$ is its differential. See more details in [3].

We consider the Lebesgue measure $dx = dX^1 \dots dX^N$. This gives to G a (bi-invariant) Haar measure (the lift of the Lebesgue measure on \mathfrak{L}) on G , that we still denote with dx . In the following, we deal with L^p spaces on G with respect to this measure. We recall that the Haar measure is invariant in the sense that it satisfies $\int f(x) dx = \int f(y \cdot x) dx$ for any $y \in G$.

Example 2.1. It is interesting to recall that all the definitions given above are generalizations of the standard properties of the Euclidean space, i.e. \mathbb{R}^n endowed with the vector sum $(x^1, \dots, x^N) + (y^1, \dots, y^N) = (x^1 + y^1, \dots, x^N + y^N)$. Indeed, \mathbb{R}^n itself is a Lie algebra with trivial brackets $[A, B] = 0$. It is realized as a matrix algebra by writing each element (x^1, \dots, x^N) as the matrix of dimension $N + 1$

$$A = \begin{pmatrix} 0 & \dots & 0 & x^1 \\ 0 & \dots & 0 & x^2 \\ \vdots & \dots & \vdots & \vdots \\ 0 & \dots & 0 & x^N \\ 0 & \dots & 0 & 0 \end{pmatrix}.$$

The exponential G satisfies $e^A = I + A$, and indeed in this case we have $e^A e^B = e^{A+B}$, thus G is isomorphic to the Euclidean group. The Haar measure is the standard Lebesgue measure on \mathbb{R}^N , and the invariance of the Haar measure is indeed the invariance of the Lebesgue measure with respect to translations.

Example 2.2. The simplest (nontrivial) example of nilpotent Lie algebra and Lie group is given by the **Heisenberg group** H_2 , see e.g. [4, Chap. 1]. Choose the Lie algebra \mathfrak{L} as the vector space generated by $\{l^1, l^2, l^3\}$ where

$$l^1 = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad l^2 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{pmatrix} \quad l^3 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}. \quad (3)$$

¹ The notation X^i with $i \in \{1, \dots, N\}$ denotes, all along the paper, a vector of the basis of the Lie algebra.

They satisfy the following commutation rules: $[l^1, l^2] = l^3$, $[l^1, l^3] = [l^2, l^3] = 0$. Hence, \mathfrak{L} is a nilpotent Lie algebra. Computing the matrix exponential, one can write the Heisenberg group as the 3D group of matrices

$$H_2 = \left\{ \begin{pmatrix} 1 & x & z + \frac{1}{2}xy \\ 0 & 1 & y \\ 0 & 0 & 1 \end{pmatrix} \mid x, y, z \in \mathbb{R} \right\}$$

endowed with the standard matrix product. Left-invariant vector fields are of the form $X(g) = gl$, where $g \in H_2$ and $l \in \mathfrak{L}$.

One can write H_2 as \mathbb{R}^3 , via the identification

$$(x, y, z) \sim \begin{pmatrix} 1 & x & z + \frac{1}{2}xy \\ 0 & 1 & y \\ 0 & 0 & 1 \end{pmatrix}.$$

In this case, the group law is

$$(x_1, y_1, z_1) \cdot (x_2, y_2, z_2) = \left(x_1 + x_2, y_1 + y_2, z_1 + z_2 + \frac{1}{2}(x_1y_2 - x_2y_1) \right).$$

The corresponding left-invariant vector fields are generated by

$$X^1 = \partial_x - \frac{y}{2}\partial_z, \quad X^2 = \partial_y + \frac{x}{2}\partial_z, \quad X^3 = \partial_z.$$

A family of dilations on \mathfrak{L} is a one-parameter family $\{\Gamma_r \mid 0 < r < \infty\}$ of automorphisms of \mathfrak{L} of the form $\Gamma_r = \exp(A \log r)$, where A is a diagonalizable linear transformation of \mathfrak{L} with positive eigenvalues. Dilations of \mathfrak{L} lift to dilations of G , that we still denote by Γ_r . The constant $Q = \text{Tr}(A)$ is called *homogeneous dimension* of G . A measurable function f on G is called *homogeneous of degree* λ if $f \circ \Gamma_r = r^\lambda f$ for all $r > 0$. A differential operator D is called *homogeneous of degree* λ if $D(u \circ \Gamma_r) = r^\lambda (Du) \circ \Gamma_r$ for all $u \in C_c^\infty(G)$ and $r > 0$. We have $d(\Gamma_r x) = r^Q dx$. For proofs and further results, see [9].

We focus on a class of nilpotent Lie algebras, the **stratified Lie algebras**. A stratification of a Lie algebra \mathfrak{L} is a decomposition $\mathfrak{L} = \bigoplus_{i=1}^s \mathfrak{L}_i$ satisfying $[\mathfrak{L}_1, \mathfrak{L}_i] = \mathfrak{L}_{i+1}$ for $1 \leq i < s$ and $[\mathfrak{L}_1, \mathfrak{L}_s] = 0$. If \mathfrak{L} is stratified, then it admits a family of dilations

$$\Gamma_r(X_1 + \cdots + X_s) = rX_1 + r^2X_2 + \cdots + r^sX_s,$$

where² $X_i \in \mathfrak{L}_i$. A **stratified group** G is the exponential of a stratified Lie algebra. In this case we have $Q = \sum_{i=1}^s i \dim(\mathfrak{L}_i)$. Given a vector field X and considering it as a differential operator, it is homogeneous of degree i if and only if $X \in \mathfrak{L}_i$. From now on, we denote s the step of nilpotency of G and w^i the degree of homogeneity of X^i . We also denote $n = \dim(\mathfrak{L}_1)$.

Example 2.3. Once again, the Euclidean group is a particular case of the stratified algebras, in which the dilation is the standard ‘‘isotropic’’ dilation $\Gamma_r(x^1, \dots, x^N) = (rx^1, \dots, rx^N)$. The homogeneous dimension coincides with N . Elements of \mathfrak{L} are considered as derivatives simply via the identification

² The notation X_i with $i \in \{1, \dots, s\}$ denotes, all along the paper, a generic vector such that $X_i \in \mathfrak{L}_i$.

$x^i \rightarrow \partial_i$, and they are indeed homogenous of order 1 with respect to the dilation.

The basic idea of Carnot groups is, more in general, that dilations can have different degrees of homogeneity with respect to different axes, that corresponds to “anisotropic” dilations.

Example 2.4. The Heisenberg group is a stratified Lie group, choosing $\mathfrak{L}_1 = \text{span}\{X^1, X^2\}$ and $\mathfrak{L}_2 = \text{span}\{X^3\}$. Its homogeneous dimension is 4.

Given a stratified Lie group G , we define a quadratic form $\langle \cdot, \cdot \rangle_S$ on \mathfrak{L}_1 , a structure that is called a **Carnot group**. Carnot groups are naturally endowed with a metric structure, the so called Carnot–Carathéodory metric. We don’t use this metric in this article. For further references, see e.g. [4].

In the following we use another norm on a Carnot group. Define a quadratic form $\langle \cdot, \cdot \rangle_R$ on the whole Lie algebra \mathfrak{L} such that it coincides with $\langle \cdot, \cdot \rangle_S$ on \mathfrak{L}_1 and such that the \mathfrak{L}_i are orthogonal with respect to this quadratic form. We introduce the corresponding norm $|v|_R := \langle v, v \rangle_R$ on \mathfrak{L} and define the following norm on G , that is homogeneous of degree 1.

$$\|\exp(X_1 + X_2 + \dots + X_s)\|_R := \left(\sum_{i=1}^s |X_i|_R^{\frac{2(s!)}{i}} \right)^{\frac{1}{2(s!)}}$$

Example 2.5. The Heisenberg group is a Carnot group when endowed with a quadratic form on \mathfrak{L}_1 . The standard choice is given by choosing X^1, X^2 orthonormal. We also endow the whole \mathfrak{L} with a quadratic form, by choosing X^1, X^2, X^3 orthonormal. Thus $\|(x, y, z)\|_R = \sqrt[4]{|x|^4 + |y|^4 + |z|^2}$.

It is interesting to remark that the quadratic form $\langle \cdot, \cdot \rangle_R$ is not intrinsic on the Carnot group $(G, \langle \cdot, \cdot \rangle_S)$. Nevertheless, all the possible norms are equivalent, i.e. given two quadratic forms $\langle \cdot, \cdot \rangle_{R_1}$ and $\langle \cdot, \cdot \rangle_{R_2}$ on \mathfrak{L} , there exist $c, C > 0$ such that $c \langle v, v \rangle_{R_1} \leq \langle v, v \rangle_{R_2} \leq C \langle v, v \rangle_{R_1}$. We assume that the vectors $\{X^1, \dots, X^n\}$ are orthonormal with respect to $\langle \cdot, \cdot \rangle_S$ and that $\{X^1, \dots, X^N\}$ are orthonormal with respect to $\langle \cdot, \cdot \rangle_R$.

2.1. Hypoelliptic heat equations on Carnot groups

As already stated, one of the main features of Carnot groups is that they are naturally endowed with a Laplacian operator. Given $\{X^1, \dots, X^n\}$ an orthonormal basis of \mathfrak{L}_1 with respect to $\langle \cdot, \cdot \rangle_S$, the intrinsic Laplacian for G is the sum of squares

$$\Delta_{sr} := \sum_{i=1}^n X_i^2.$$

Here, we are interested in the **hypoelliptic heat operator** $\partial_t - \Delta_{sr}$.

The Laplacian and the heat operator are intrinsic in the sense that they are intrinsically defined from the structure of the Carnot group. In particular, one can find connections between the solution of the hypoelliptic heat equation and the intrinsic (sub-Riemannian) distance defined on Carnot group. This interesting topic is outside the goals of this paper: more informations and

interesting estimates can be found in [1, 6, 8, 15]. Other results for the hypoelliptic heat operator defined on other spaces than Carnot groups can be found in [2, 14].

We recall that both the Laplacian and the heat operator are hypoelliptic, according to the following definition (see e.g. [12]).

Definition 2.6. An operator L is hypoelliptic when, given $U \subset \mathbb{R}^N$ and $\phi : U \rightarrow \mathbb{C}$ such that $L\phi \in C^\infty$, then ϕ is C^∞ .

We are interested in the **hypoelliptic heat equation** on Carnot groups. We recall some properties.

Theorem 2.7. ([9] (2.1), (2.11), (3.1)) *Let G be a Carnot group of homogeneous dimension $Q > 2$. Then the heat operator $\partial_t - \Delta_{sr}$ is homogeneous of degree 2. There is a unique heat kernel*

$$P : G \times (0, +\infty) \rightarrow \mathbb{R}$$

$$(x, t) \mapsto P_t(x)$$

of type 2 such that $P_0(\cdot) = \delta_0$. It satisfies $P_t(x) \geq 0$, $\int_G P_t(x) dx = 1$ for all t and $P_{r^{2t}}(\Gamma_r(x)) = r^{-Q} P_t(x)$.

For the following, we need estimates about the large time behavior of the heat kernel P_t and of its derivatives (of any order) with respect to left-invariant vector fields. These results are proved in [9].

Corollary 2.8. ([9] (3.4)–(3.6)) *The kernel $P_t(x)$ is C^∞ on $G \times (0, +\infty)$. For each $t_0 > 0$ and positive integer k there is a constant $C_{t_0, k}$ such that, for all $\|x\|_R \geq 1$ and $t \leq t_0$, we have*

$$|P_t(x)| \leq C_{t_0, k} \|x\|_R^{-k}. \tag{4}$$

Let D be any left-invariant differential operator on G . Then the same estimate (4) holds for DP_t .

We now give estimates for the L^p norms of the heat kernel and its derivatives.

Lemma 2.9. *Let G be a Carnot group of homogeneous dimension $Q > 2$, P_t the corresponding heat kernel, with $t > 0$. Given $1 \leq q \leq \infty$, we have $\|P_1\|_q < \infty$ and $\|P_t\|_q = t^{-Q/(2q^*)} \|P_1\|_q$, with $\frac{1}{q} + \frac{1}{q^*} = 1$.*

Given an homogeneous left-invariant differential operator D of degree λ , we have $\|DP_1\|_q < \infty$ and $\|DP_t\|_q = t^{-Q/(2q^) - \lambda/2} \|DP_1\|_q$.*

Proof. For $1 < q < \infty$ we have $\|P_t\|_q^q = \int_G |P_t(x)|^q dx = \int_G r^{Qq} |P_{r^{2t}}(\Gamma_r(x))|^q dx = r^{Qq} \int_G |P_{r^{2t}}(y)|^q r^{-Q} dy = r^{Q(q-1)} \|P_{r^{2t}}\|_q^q$. The proof is given by choosing $r = t^{-\frac{1}{2}}$. Remark that $\|P_1\|_q$ is finite due to (4). The proof for $q = 1$ and $q = \infty$ is equivalent. The proofs for DP_t are identical, recalling $DP_{r^{2t}}(\Gamma_r x) = r^{-Q-\lambda} DP_t(t, x)$. \square

2.2. Solutions of the initial value problem

Given the fundamental solution P_t for the hypoelliptic heat equation, we want to find the solution of the initial value problem (2) for a given f_0 . The solution is given by the **convolution** $f_t = f_0 * P_t$, with the convolution defined on Lie groups as follows.

Definition 2.10. Let (G, \cdot) be a Lie group and dg its left-invariant Haar measure. The convolution is

$$(f_1 * f_2)(g) = \int_G f_1(h) f_2(h^{-1} \cdot g) dh.$$

Remark that the product and the inverse are given by the group operation on G . The convolution on Lie groups satisfies an invariance property with respect to left-invariant differential operators. More precisely, given a left-invariant vector field X , we have

$$X(f_1) * f_2 = X(f_1 * f_2) = f_1 * X(f_2). \quad (5)$$

Remark 2.11. The standard definition of convolution on \mathbb{R}^N is given by considering it as the Euclidean space. In this case, left-invariant vector fields are differential operators with constant coefficient.

A useful property is the Young's inequality

Theorem 2.12. ([11] (20.14–18)) *Let $f_1 \in L^1(G)$ and $f_2 \in L^p(G)$. Then $f_1 * f_2 \in L^p(G)$ and*

$$\|f_1 * f_2\|_p \leq \|f_1\|_1 \|f_2\|_p.$$

3. Functions decomposition on Carnot groups

In this section, we generalize to Carnot groups a decomposition of functions based on Dirac delta and its derivatives, first stated in [7] for functions defined on \mathbb{R}^N . We first recall that, given $F = (F_1, \dots, F_N)$, the intrinsic divergence³ is $\operatorname{div} F = \sum_{i=1}^N X^i(F_i)$.

We give a first decomposition of functions, based on Dirac delta only. The result is not interesting in itself, since it is weaker than [7, Theorem 1], but its proof gives an example of the method used in the following to deal with Carnot groups.

Proposition 3.1. *Let f be a function of G and $p \in [1, \infty]$.*

- *If $p < Q/(Q - 1)$, $f \in L^1(G)$ and $|x|_R f \in L^p(G)$, then it exists $F \in (L^p(G))^N$ such that $f = (\int f) \delta_0 + \operatorname{div} F$.*
- *If $p > Q/(Q - s)$ and $|x|_R f \in L^p(G)$, then it exists $F \in (L^p(G))^N$ such that $f = \operatorname{div} F$.*

We also have that $\|F_j\|_p \leq C_{p,N} \| |x|_R f \|_p$ for each $j = 1, \dots, N$.

³ Also in this case, the divergence is intrinsic in the sense that is intrinsically defined by the structure of the Carnot group. For more details, see e.g [1].

Proof. The proof is a slight generalization of the proof given in [7]. Fix $x = (x_1, \dots, x_N) \in G$ and take a test function $\phi \in \mathfrak{D}(G)$. For the first case, we apply the fundamental theorem of integral calculus for the function of $r \in [0, 1]$ defined by $\tilde{\phi}(r) := \phi(\Gamma_r(x))$ finding

$$\begin{aligned} \phi(x) - \phi(0) &= \int_0^1 \frac{\partial \tilde{\phi}}{\partial r} dr = \int_0^1 \sum_{i=1}^N X^i(\phi)|_{\Gamma_r(x)} \frac{\partial(r^{w^i} x_i)}{\partial r} dr \\ &= \int_0^1 \sum_{i=1}^N w^i r^{w^i-1} x_i X^i(\phi)|_{\Gamma_r(x)} dr. \end{aligned}$$

We now apply $f - (\int f) \delta_0$ to $\phi(x)$ and have

$$\begin{aligned} \int_G \left(f - \left(\int f \right) \delta_0 \right) \phi dx &= \int_G f(x) \int_0^1 \sum_{i=1}^N w^i x_i r^{w^i-1} X^i(\phi)|_{\Gamma_r(x)} dr dx \\ &= - \int_G \phi(x) \sum_{i=1}^N X^i \left(w^i x_i \int_0^1 \frac{f(\Gamma_{r^{-1}}(x))}{r^{Q+1}} dr \right) dx. \end{aligned}$$

The last identity is given by integration by parts, since the X^i are skew-adjoint with respect to $dx = dX^1 \dots dX^N$, see e.g. [1, Pr. 20]. We also apply change of variables $\Gamma_r(x) \mapsto x'$. Choose $F_i(x) := w^i x_i \int_0^1 \frac{f(\Gamma_{r^{-1}}(x))}{r^{Q+1}} dr$. In this case we have

$$\|F_i\|_p^p = |w^i|^p \cdot \| |x|_R f(x) \|_p^p \cdot \left| \int_0^1 r^{\frac{Q}{p} + w^i - Q - 1} dr \right|^p. \quad (6)$$

It is finite when $p < \frac{Q}{Q-w^i}$. Since $w^i \geq 1$, it is finite when $p < \frac{Q}{Q-1}$.

The other case is similar, given the formula

$$\phi(x) = - \int_0^1 \sum_{i=1}^N w^i r^{-1-w^i} x_i X^i(\phi(\Gamma_{r^{-1}}(x))) dr,$$

from which we find

$$F_i(x) := w^i x_i \int_0^1 r^{Q-1} f(\Gamma_r(x)) dr. \quad (7)$$

In this case, we have

$$\|F_i\|_p^p = |w^i|^p \cdot \| |x|_R f(x) \|_p^p \cdot \left| \int_0^1 r^{-\frac{Q}{p} + Q - 1 - w^i} dr \right|^p.$$

It is finite when $p > \frac{Q}{Q-w^i}$. Since $w_i \leq s$, it is finite when $p > \frac{Q}{Q-s}$. \square

Remark 3.2. One can check that all the formulas reduce to the results of [7] when $G = (R^N, +)$, recalling that $s = 1$, $w_i \equiv 1$, $\Gamma_r(x) = rx$ and $Q = N$. The main difference with [7] is the set of L^p spaces in which none of the decompositions holds, that is $p \in \left[\frac{Q}{Q-1}, \frac{Q}{Q-s} \right]$.

Remark 3.3. One can ask if a similar decomposition can hold for $p \in \left[\frac{Q}{Q-1}, \frac{Q}{Q-s} \right]$. It is clearly possible, since one can take results of [7, Theorem 1], provided $p \neq \frac{N}{N-1}$. Remark that $\frac{N}{N-1} \in \left[\frac{Q}{Q-1}, \frac{Q}{Q-s} \right]$.

If one looks for a formula $f = (\int f) \delta_0 + g$ with g given by a differential operator that is not the divergence, it is certainly possible to find more complicated decomposition. We use this idea in the following, that gives operators in which vector fields X^i play different roles depending on their weight w_i .

As stated above, a higher order decomposition can be defined using the derivatives of the Dirac delta as basis functions, following [7, Theorem 2]. In this case, it is necessary to compute high order derivatives of $\tilde{\phi}(s)$. For this reason, we only state the first order decomposition (its meaning will be clear in the following). For higher order decomposition one has to apply the same technique with longer computations.

Theorem 3.4. *Let f be a function of G and $p \in [1, \infty]$.*

- *If $p < \frac{Q}{Q-1}$, $f \in L^1(G, 1 + |x|_R)$ and $|x|_R f, |x|_R^2 f \in L^p(G)$, then*

$$f = \left(\int_G f \right) \delta_0 - \sum_{i=1}^n \left(\int f x_i \right) (X_i \delta_0) + \sum_{i,j=1}^N X^i X^j (F_{ij}) + \sum_{i=n+1}^N X^i (F_i).$$

- *If $p > \frac{Q}{Q-2s}$ and $|x|_R f, |x|_R^2 f \in L^p(G)$, then*

$$f = \sum_{i,j=1}^N X^i X^j (F_{ij}) + \sum_{i=n+1}^N X^i (F_i).$$

In both cases, we have $\|F_{ij}\|_p \leq C \| |x|_R^2 f \|_p$ and $\|F_i\|_p \leq C \| |x|_R f \|_p$.

Proof. Fix $x = (x_1, \dots, x_N) \in G$. Given a test function $\phi \in \mathfrak{D}(G)$, consider the function of $r \in [0, 1]$ defined by $\tilde{\phi}(r) := \phi(\Gamma_r(x))$. The Taylor polynomial $P_k \tilde{\phi}(r)$ of order k near 0 for $\tilde{\phi}$ satisfies

$$\tilde{\phi}(1) - P_k \tilde{\phi}(1) = \int_0^1 \frac{\tilde{\phi}^{(k+1)}(r)}{k!} (1-r)^k dr.$$

We also have

$$\begin{aligned} \tilde{\phi}(1) &= (-1)^{k+1} \int_1^\infty \frac{\tilde{\phi}^{(k+1)}(r)}{k!} (r-1)^k dr \\ &= (-1)^{k+1} \int_0^1 \frac{\tilde{\phi}^{(k+1)}\left(\frac{1}{t}\right)}{k!} (1-t)^k \frac{dt}{t^{k+2}}, \end{aligned} \tag{8}$$

that is a particular case⁴ of

$$\tilde{\phi}(r) = (-1)^{k+1} \int_0^1 \frac{\tilde{\phi}^{(k+1)}\left(\frac{r}{t}\right)}{k!} \left(\frac{r}{t}\right)^{k+1} (1-t)^k \frac{dt}{t}.$$

⁴ Remark a misprint in [7, Eq. (6)], where $(-1)^k$ must be replaced with $(-1)^{k+1}$.

We merge the two in this more general formula

$$\begin{aligned}\tilde{\phi}(1) - P_{k-j}\tilde{\phi}(1) &= \int_0^1 \frac{\tilde{\phi}^{(k-j+1)}(r)}{(k-j)!} (1-r)^{k-j} dr \\ &= (-1)^j \int_0^1 dr \int_0^1 \frac{dt}{t} \left(\frac{r}{t}\right)^j (1-t)^{j-1} (1-r)^{k-j} \frac{\tilde{\phi}^{(k+1)}\left(\frac{r}{t}\right)}{(k-j)!(j-1)!}.\end{aligned}\quad (9)$$

We now restrict ourselves to the case $k = 1$. We compute $\tilde{\phi}'$ and $\tilde{\phi}''$:

$$\begin{aligned}\tilde{\phi}'(r) &= \sum_{i=1}^N X^i(\phi)|_{\Gamma_r(x)} w^i r^{w^i-1} x_i, \\ \tilde{\phi}''(r) &= \sum_{i,j=1}^N X^i X^j(\phi)|_{\Gamma_r(x)} w^i w^j r^{w^i+w^j-2} x_i x_j \\ &\quad + \sum_{i=n+1}^N X^i(\phi)|_{\Gamma_r(x)} w^i (w^i - 1) r^{w^i-2} x_i.\end{aligned}\quad (10)$$

Remark that the last term in the r.h.s. of (10) contains first order derivatives with respect to vector fields whose degree of homogeneity is ≥ 2 . We compute the Taylor formulas for $\tilde{\phi}$:

$$P_0\tilde{\phi}(1) = \phi(0), \quad P_1\tilde{\phi}(1) = \phi(0) + \sum_{i=1}^n X^i(\phi)|_0 x_i.$$

Remark that $P_1\tilde{\phi}$ contains derivatives in \mathfrak{L}_1 only.

We now apply $f - ((\int_G f)\delta_0 - \sum_{i=1}^n (\int f x_i)(X_i \delta_0))$ to ϕ , that gives

$$\begin{aligned}&\int \left(f(x)\phi(x) - f(x)\phi(0) - \sum_{i=1}^n f(x)x_i X^i(\phi)|_0 \right) dx \\ &= \int f(x) \int_0^1 \frac{\tilde{\phi}''(r)}{k!} (1-r)^k dr dx \\ &= \int \frac{f(x)}{k!} \left(\int_0^1 \sum_{i,j=1}^N X^i X^j(\phi)|_{\Gamma_r(x)} w^i w^j r^{w^i+w^j-2} x_i x_j \right. \\ &\quad \left. + \sum_{i=n+1}^N X^i(\phi)|_{\Gamma_r(x)} w^i (w^i - 1) r^{w^i-2} x_i \right) (1-r) dr dx \\ &= \int_G \phi(x) \left(\sum_{i,j=1}^N \int_0^1 w^i w^j X^i X^j (f(\Gamma_{r-1}(x))x_i x_j) \frac{1-r}{r^{Q+2}} \right. \\ &\quad \left. - \sum_{i=n+1}^N \int_0^1 w^i (w^i - 1) X^i (f(\Gamma_{r-1}(x))x_i) \frac{1-r}{r^{Q+2}} \right) dr dx.\end{aligned}$$

Choose $F_{ij} = F_{ji} := \int_0^1 w^i w^j f(\Gamma_{r^{-1}}(x)) x_i x_j \frac{1-r}{rQ+2} dr$, that satisfies

$$\|F_{ij}\|_p^p = |w^i w^j|^p \int |f(x) x_i x_j|^p dx \left| \int_0^1 r^{w^i + w^j - 2 - Q + \frac{Q}{p}} (1-r) dr \right|^p,$$

that is finite when $f(x) x_i x_j \in L^p(G)$ and $w^i + w^j - 2 - Q + \frac{Q}{p} > -1$. Since $w_i \geq 1$, it is finite when $p < \frac{Q}{Q-1}$. Similarly, for $i \geq n+1$ choose $F_i(x) := \int_0^1 w^i (w^i - 1) f(\Gamma_{r^{-1}}(x)) x_i \frac{1-r}{rQ+2} dr$, that satisfies

$$\|F_i\|_p^p = |w^i (w^i - 1)|^p \int |f(x) x_i|^p dx \left| \int_0^1 r^{w^i - 2 - Q - \frac{Q}{p}} (1-r) dr \right|^p.$$

It is finite when $f(x) x_i \in L^p(G)$ and $w^i - 2 - Q + \frac{Q}{p} > -1$. Since $w^i \geq 2$, it is finite when $p < \frac{Q}{Q-1}$.

The second case is similar, using (8) and (10). Second order terms give condition $p > \frac{Q}{Q-2s}$, while first order terms give $p > \frac{Q}{Q-s}$. \square

Remark 3.5. Observe that the decomposition given above is not a particular case of results of [7]. Indeed, there are two main differences with respect to similar results stated in \mathbb{R}^N . In the first formula, one can observe that the first-order term $-\sum_{i=1}^n (\int f x_i)(X_i \delta_0)$ contains derivatives in \mathfrak{L}_1 only, and not in the whole algebra \mathfrak{L} as it is in \mathbb{R}^N . The second difference is related to the remainder $\sum_{i,j=1}^N X^i X^j (F_{ij}) + \sum_{i=n+1}^N X^i (F_i)$, that contains both second order derivatives with respect to all vector fields, and first order derivatives of vector fields outside \mathfrak{L}_1 . Both results can be generalized by defining polynomials on Carnot groups, with the degree being indeed the degree of homogeneity. For details, see [4].

Remark 3.6. Comments given in Remark 3.2 apply to this case too. The computation of higher order derivatives is very simple for isotropic dilations, as shown both in [4] and [7]. But also in this case, we want to highlight the different weights of the vector fields X^i .

Remark 3.7. Taylor formula for ϕ is indeed a particular case of formulas in [4, Sect. 20.3.2] and [5]. We have also simplified the notation, using coordinate x_i instead of exponentiation and logarithm.

Remark 3.8. Following [7], one could ask if a decomposition of the kind

$$f = \left(\int_G f \right) \delta_0 + \sum_{i,j=1}^N X^i X^j (F_{ij}) + \sum_{i=n+1}^N X^i (F_i) \quad (11)$$

can be found. We apply $f - (\int f) \delta_0$ to ϕ and use (9) with $k = j = 1$. We find formula (11) with

$$F_{ij}(x) := -w^i w^j x_i x_j \int_0^1 \int_0^1 f \left(\Gamma_{\left(\frac{r}{t}\right)^{-1}}(x) \right) \left(\frac{r}{t} \right)^{-Q-1} dr \frac{dt}{t}$$

and

$$F_i := w^i (w^i - 1) x_i \int_0^1 \int_0^1 f \left(\Gamma_{\left(\frac{r}{t}\right)^{-1}}(x) \right) \left(\frac{r}{t} \right)^{-Q-1} dr \frac{dt}{t}.$$

We have that $F_{ij} \in L^p(G)$ when $f(x)x_ix_j \in L^p(G)$ and $-1 < w^i + w^j - 1 - Q + \frac{Q}{p} < 0$. This formula must hold for all w_i, w_j . In particular, $p > \frac{Q}{Q-2s+1}$ and $p < \frac{Q}{Q-2}$. This implies that $s < 2$, i.e. $s = 1$. It means that G is a 1-step Carnot group, i.e. a standard Euclidean space. Thus, the decomposition coincides with the one given in [7, Theorem 2(a)] with $j = k = 1$. In other terms, this decomposition only holds for Euclidean spaces.

4. Large time behavior for the Cauchy problem with hypoelliptic heat diffusion

We present here the large time behavior of the solutions of (2).

Theorem 4.1. *Let $f_0 \in L^1(G)$ satisfying $|x|_R f \in L^p(G)$ with $p < \frac{Q}{Q-1}$. Let $f(t)$ be the solution at time $t > 0$ of (2). Then, given $p \leq q \leq \infty$, there exists C , not depending on t and f_0 , such that*

$$\left\| f(t) - \left(\int f_0 \right) P_t \right\|_q \leq C t^{-Q((1/p)-(1/q))/2-1/2} \| |x|_R f_0 \|_p.$$

Proof. We have $f_0 = (\int f_0) \delta_0 + \operatorname{div} F$, where $F_i \in L^p(G)$. Recall that $f(t) = f_0 * P_t$, thus

$$\begin{aligned} \left\| f(t) - \left(\int f_0 \right) P_t \right\|_q &= \left\| \sum_{i=1}^N X^i F_i * P_t \right\|_q = \left\| \sum_{i=1}^N F_i * X^i P_t \right\|_q \\ &\leq \sum_{i=1}^N \| F_i \|_p \| X^i P_t \|_r, \end{aligned}$$

where $\frac{1}{p} + \frac{1}{r} = \frac{1}{q} + 1$. Recall that $\| F_i \|_p \leq C_{p,G} \| |x|_R f_0 \|_p$ and that each X^i is homogeneous of order greater than or equal to 1. Hence $\| f(t) - (\int f_0) P_t \|_q \leq \sum_{i=1}^N C_{p,G} \| |x|_R f \|_p t^{-Q/(2r^*)-1/2} \| X^i P_1 \|_r$, with $\frac{1}{r^*} = \frac{1}{p} - \frac{1}{q}$. All the terms $\| X^i P_1 \|_r$ are finite (Lemma 2.9). Consequently, the statement of the theorem holds with $C := \sum_{i=1}^N C_{p,G} \| X^i P_1 \|_r$. \square

We prove now a finer estimate, using Theorem 3.4.

Theorem 4.2. *Let $f_0 \in L^1(G, 1 + |x|_R)$ such that $|x|_R f, |x|_R^2 f \in L^p(G)$ with a given $p < \frac{Q}{Q-1}$. Let $f(t)$ be the solution at time $t > 0$ of (2). Then, given $p \leq q \leq \infty$, there exists C , not depending on t and f_0 , satisfying*

$$\begin{aligned} \left\| f(t) - \left(\left(\int f_0 \right) P_t + \sum_{i=1}^n \left(\int f_0 x_i \right) X^i(P_t) \right) \right\|_q \\ \leq C t^{-Q((1/p)-(1/q))/2-1} \| |x|^{k+1} f_0 \|_p. \end{aligned}$$

Proof. Using Theorem 3.4, we have

$$\begin{aligned}
 & \left\| f(t) - \left(\left(\int f_0 \right) P_t + \sum_{i=1}^n \left(\int f_0 x_i \right) X^i(P_t) \right) \right\|_q \\
 &= \left\| \sum_{i,j=1}^N X^i X^j(F_{ij}) P_t + \sum_{i=n+1}^N X^i(F_i) P_t \right\|_q \\
 &\leq \sum_{i,j=1}^N \|F_{ij}\|_p \|X^i X^j(P_t)\|_r + \sum_{i=n+1}^N \|F_i\|_p \|X^i(P_t)\|_r \\
 &\leq \sum_{i,j=1}^N \| |x|_R^2 f_0 \|_p t^{-Q/(2r^*)-1} \|X^i X^j P_1\|_r \\
 &\quad + \sum_{i=n+1}^N \| |x|_R f_0 \|_p t^{-Q/(2r^*)-1} \|X^i P_1\|_r,
 \end{aligned}$$

where $\frac{1}{p} + \frac{1}{r} = \frac{1}{q} + 1$, thus $\frac{1}{r^*} = \frac{1}{p} - \frac{1}{q}$. We have used here Lemma 2.9, recalling that derivatives $X^i X^j$ are homogeneous of order at least 2, and that the same holds for X^i with $i = n + 1, \dots, N$. \square

This expression shows that, for the large time behavior, the main contribution is given by P_t , its derivatives $X^1 P_t, \dots, X^n P_t$ in \mathfrak{L}_1 (case $w^i = 1$) and, in general, by terms with small degree of homogeneity. This is a direct consequence of the anisotropy of the diffusion on stratified Lie groups.

Higher-order estimations can be easily proved with the same techniques, provided higher-order decomposition of f_0 of the kind of Theorem 3.4.

4.1. The Heisenberg group

We now apply the results presented before to the **Heisenberg group**, defined in Examples 2.2, 2.4, 2.5. The hypoelliptic Laplacian Δ_{sr} is the sum of squares

$$\Delta_{sr}\phi = \left((X^1)^2 + (X^2)^2 \right) \phi.$$

The explicit expression of P_t has been first computed in [10, 13]. We give here the expression with respect to the definition of H_2 given above, see [1]:

$$P_t(x, y, z) = \frac{1}{(2\pi t)^2} \int_{\mathbb{R}} \frac{2\tau}{\sinh(2\tau)} \exp\left(-\frac{\tau(x^2 + y^2)}{2t \tanh(2\tau)}\right) \cos\left(\frac{z\tau}{t}\right)$$

We consider a given function $f_0 \in L^1(H_2, 1 + |x|_R)$ such that $|x|_R f_0, \|x\|_R^2 f_0 \in L^p$ for a fixed $p < \frac{4}{3}$. We choose $p = 1$. Thus f_0 satisfies the hypotheses of Theorem 3.4. We compute the momenta A_i defined as follows

$$\begin{aligned}
 A_0 &:= \int_{H_2} |f_0(x, y, z)| dx dy dz, \\
 A_1 &:= \int_{H_2} |f_0(x, y, z)x| dx dy dz, \quad A_2 := \int_{H_2} |f_0(x, y, z)y| dx dy dz.
 \end{aligned}$$

We now apply Theorem 4.2 with $q = \infty$, that provides the large time expression in $L^\infty(H_2)$ for the solution of (2) on H_2 :

$$f(t) = A_0 P_t + A_1 X^1(P_t) + A_2 X^2(P_t) + O(t^{-2/p-1}).$$

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Large time for heat on Carnot groups

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